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# AN OBSTRUCTION FOR THE MEAN CURVATURE OF A CONFORMAL IMMERSION $S^n \rightarrow \mathbb{R}^{n+1}$

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**ABSTRACT.** We prove a Pohozaev type identity for non-linear eigenvalue equations of the Dirac operator on Riemannian spin manifolds with boundary. As an application, we obtain that the mean curvature  $H$  of a conformal immersion  $S^n \rightarrow \mathbb{R}^{n+1}$  satisfies  $\int \partial_X H = 0$  where  $X$  is a conformal vector field on  $S^n$  and where the integration is carried out with respect to the Euclidean volume measure of the image. This identity is analogous to the Kazdan-Warner obstruction that appears in the problem of prescribing the scalar curvature on  $S^n$  inside the standard conformal class.

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Let  $(M, g)$  be a compact Riemannian manifold with a conformal vector field  $X$ . Given a function  $s$  on  $M$ , then it is a classical question to ask whether  $s$  is the scalar curvature of a metric  $\tilde{g}$  conformal to  $g$ . The determination of the set of all such functions  $s$  is still open, although several partial results are known, in particular there are necessary conditions that  $s$  has to satisfy in order to be a scalar curvature.

On the one hand there are topological obstructions. If for example  $M$  is spin and has non-vanishing  $\hat{A}$  genus, then the scalar curvature of any metric on  $M$  has either to be negative somewhere or the Ricci curvature vanishes everywhere on  $M$ .

However, if one fixes the conformal class  $[g]$  as described above, there are further obstructions that arise from conformal vector fields. For example if  $M$  is  $S^n$  with the standard conformal structure, Kazdan and Warner [KW75] derived a famous obstruction. A slightly stronger version of this obstruction due to Bourguignon and Ezin [BE87] is described in the following theorem.

**Theorem 1.** *Let  $X$  be a conformal vector field on the compact manifold  $(M, g)$ . If  $s$  is the scalar curvature of a metric  $\tilde{g} = u^{4/(n-2)}g$ , then*

$$\int_M \partial_X s \, dv_{\tilde{g}} = 0$$

where  $dv_{\tilde{g}} = u^{\frac{2n}{n-2}} dv_g$  is the volume measure associated to  $\tilde{g}$ .

Tightly related to the Kazdan-Warner obstruction is the Pohozaev identity. Let  $\Omega$  be a star-shaped open set of  $\mathbb{R}^n$  ( $n \in \mathbb{N}$ ) with smooth boundary. We denote by  $\Delta = -\sum_{i=1}^n \partial_{ii}$  the Laplacian on  $\mathbb{R}^n$ . Let  $u \in C^2(\bar{\Omega})$  be a positive solution of  $\Delta u = u^{p-1}$  on  $\Omega$  with  $u|_{\partial\Omega} \equiv 0$ . The vector field  $X = \sum_{i=1}^n x^i \partial_i$  is conformal. If one uses similar methods as in the proof of the Kazdan-Warner obstruction, then one obtains the Pohozaev identity ([Po65]) which asserts that:

$$\left(1 - \frac{n}{2} + \frac{n}{p}\right) \int_{\Omega} u^p = \frac{1}{2} \int_{\partial\Omega} \langle \nu, X \rangle (\partial_{\nu} u)^2 \quad (1)$$

where  $\nu$  resp.  $\partial_{\nu}$  is the outer normal vector resp. the outer normal derivative on  $\partial\Omega$ . One among many important consequences of this inequality is that no non-trivial solutions exist if  $p \geq \frac{2n}{n-2}$ . Another application is an alternative proof of the Kazdan-Warner obstructions in the case that  $(M, g)$  is the sphere with the standard conformal structure [DR99].

In the present short article, we establish a similar identity for the classical Dirac operator  $D$ . We derive this identity on manifolds with boundary in order to admit future Pohozaev type applications. Then,

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we will specialize to compact manifolds without boundary, where we will derive a Kazdan-Warner type obstruction for the mean curvature of a conformal immersion  $S^2 \rightarrow \mathbb{R}^3$ .

Our main theorem is:

**Theorem 2.** *Let  $(M, g, \chi)$  be a compact Riemannian spin manifold of dimension  $n$  with boundary  $\partial M$  (possibly equal to  $\emptyset$ ) and with Dirac operator  $D$ . We assume that there exists a smooth spinor field  $\psi$  which satisfies for some  $p > 1$ ,*

$$D\psi = H|\psi|^{p-2}\psi, \quad H \in C^\infty(M). \quad (2)$$

Furthermore, we assume that  $X$  is a conformal vector field on  $M$ . Then, we have the following Pohozaev type identity:

$$\int_{\partial M} \langle \nu \cdot \mathcal{L}_X \psi, \psi \rangle = \frac{p-2}{p} \int_{\partial M} H|\psi|^p g(X, \nu) + \left( \frac{1}{n} - \frac{p-2}{p} \right) \int_M H \operatorname{div} X |\psi|^p + \frac{2}{p} \int_M (\partial_X H) |\psi|^p,$$

where  $\nu$  denotes the outward pointing normal vector along  $\partial M$ , and where  $\langle \cdot, \cdot \rangle$  denotes the real scalar product on spinors.

**Proof:** The flow associated to the conformal vector field  $X$  will be denoted as  $\alpha^t$ . If  $p$  is in the interior of  $M$ , then  $\alpha^t(p)$  exists for times  $t$  close to 0. For any  $t \in \mathbb{R}$  let  $f^t$  be the conformal scaling function of  $\alpha^t$ , i.e.  $(d\alpha^t)_p$  is  $f^t(p)$  times an isometry from  $T_p M$  to  $T_{\alpha^t(p)} M$ . Let  $\alpha_*^t : \Sigma_p M \rightarrow \Sigma_{\alpha^t(p)} M$  be the spinor identification map as constructed in [Ht74, Hi86, BG92]. In particular, this map has the pointwise properties that

$$|\alpha_*^t(\psi)| = |\psi|$$

and the following transformation formula for conformal changes of the metric. Let  $\varphi \in \Gamma(\Sigma M)$  be a spinor field. For  $t$  close to 0, we then define the map  $\alpha_\#^t : \Gamma(\Sigma M) \rightarrow \Gamma(\Sigma \tilde{M})$ ,  $\alpha_\#^t(\varphi) := \alpha_*^t \circ \varphi \circ \alpha^{-t}$ , where  $\tilde{M}$  is  $M$  without an open neighborhood of the boundary.

Then

$$D\alpha_\#^t \left( (f^t)^{-\frac{n-1}{2}} \psi \right) = \alpha_\#^t \left( (f^t)^{-\frac{n+1}{2}} D\psi \right).$$

Now we assume that  $\psi$  satisfies (2), and we obtain

$$D\alpha_\#^t \left( (f^t)^{-\frac{n-1}{2}} \psi \right) = \alpha_\#^t \left( (f^t)^{-\frac{n+1}{2}} H|\psi|^{p-2}\psi \right).$$

Deriving with respect to  $t$  at  $t = 0$  yields

$$-\frac{n-1}{2} D\beta\psi + D \frac{d}{dt} \Big|_{t=0} \alpha_\#^t \psi = -\frac{n+1}{2} H\beta|\psi|^{p-2}\psi + H|\psi|^{p-2} \frac{d}{dt} \Big|_{t=0} \alpha_\#^t \psi \quad (3)$$

$$+ (p-2)H \langle \frac{d}{dt} \Big|_{t=0} \alpha_\#^t \psi, \psi \rangle |\psi|^{p-4}\psi - (\partial_X H) |\psi|^{p-2}\psi. \quad (4)$$

where  $\beta := \frac{d}{dt} \Big|_{t=0} f^t$ . We reformulate using definition of the *Lie derivative of spinor fields in the direction  $X$*  [BG92], i.e.

$$\mathcal{L}_X(\psi) = -\frac{d}{dt} \Big|_{t=0} \alpha_\#^t(\psi). \quad (5)$$

Together with  $D\beta\psi = \beta D\psi + \nabla\beta \cdot \psi$  and (2) one then concludes that

$$\frac{n-1}{2} \nabla\beta \cdot \psi + D\mathcal{L}_X\psi = H\beta|\psi|^{p-2}\psi + H|\psi|^{p-2}\mathcal{L}_X\psi \quad (6)$$

$$+ (p-2)H \langle \mathcal{L}_X\psi, \psi \rangle |\psi|^{p-4}\psi + (\partial_X H) |\psi|^{p-2}\psi. \quad (7)$$

After multiplication with  $\psi$ , the  $\nabla\beta \cdot \psi$ -term vanishes, and we obtain

$$\langle D\mathcal{L}_X\psi, \psi \rangle = (p-1)H|\psi|^{p-2} \langle \mathcal{L}_X\psi, \psi \rangle + H\beta|\psi|^p + (\partial_X H) |\psi|^p.$$

The product rule for the Lie derivative tells us that

$$|\psi|^{p-2} \langle \mathcal{L}_X\psi, \psi \rangle = \frac{1}{2} |\psi|^{p-2} \partial_X |\psi|^2 = |\psi|^{p-1} \partial_X |\psi| = \frac{1}{p} \partial_X |\psi|^p.$$

Hence, we obtain

$$\langle D\mathcal{L}_X\psi, \psi \rangle = \frac{p-1}{p} H \partial_X |\psi|^p + H\beta |\psi|^p + (\partial_X H) |\psi|^p.$$

Strictly speaking, this equation is valid in the interior, but it extends to the boundary by continuity. Now, we integrate over  $M$ . With partial integration for the Dirac operator one obtains

$$\int_M \langle D\mathcal{L}_X\psi, \psi \rangle = \int_M \langle \mathcal{L}_X\psi, D\psi \rangle + \int_{\partial M} \langle \nu \cdot \mathcal{L}_X\psi, \psi \rangle = \int_M H \underbrace{\langle \mathcal{L}_X\psi, |\psi|^{p-2}\psi \rangle}_{=\frac{1}{p}\partial_X |\psi|^p} + \int_{\partial M} \langle \nu \cdot \mathcal{L}_X\psi, \psi \rangle.$$

This yields

$$\int_{\partial M} \langle \nu \cdot \mathcal{L}_X\psi, \psi \rangle = \frac{p-2}{p} \int_M H \partial_X |\psi|^p + \int_M H\beta |\psi|^p + \int_M (\partial_X H) |\psi|^p.$$

Using  $\operatorname{div}(H|\psi|^p X) = (\partial_X H)|\psi|^p + H\partial_X |\psi|^p + H|\psi|^p \operatorname{div} X$  and  $\operatorname{div} X = n\beta$  we obtain

$$\int_{\partial M} \langle \nu \cdot \mathcal{L}_X\psi, \psi \rangle = \frac{p-2}{p} \int_{\partial M} H |\psi|^p g(X, \nu) + \left( \frac{1}{n} - \frac{p-2}{p} \right) \int_M H \operatorname{div} X |\psi|^p + \frac{2}{p} \int_M (\partial_X H) |\psi|^p,$$

### Examples 3.

1.) Let  $\Omega$  be domain in  $\mathbb{R}^n$  with smooth boundary, let  $X = r\partial_r = \sum x^i \partial_i$ , and we will assume that  $H = \lambda$  is constant. Then  $\beta \equiv 1$  and we obtain

$$\int_{\partial\Omega} \langle \nu \cdot \mathcal{L}_X\psi, \psi \rangle = \lambda \frac{p-2}{p} \int_{\partial\Omega} \langle X, \nu \rangle |\psi|^p + \lambda \left( 1 - \frac{p-2}{p} n \right) \int_{\Omega} |\psi|^p.$$

This inequality bears many analogies to equation (1). In particular, the constant  $1 - \frac{p-2}{p}$  before the integral over  $\Omega$  vanishes if  $p$  takes the value  $p = 2n/(n-1)$ . This value plays the same role in non-linear Dirac equations as the value  $p = 2n/(n-2)$  does for the Laplace operator.

2.) If  $M$  is a closed manifold and  $X$  is a conformal vector field, then for  $p = 2n/(n-1)$  we obtain

$$\int_M (\partial_X H) |\psi|^p = 0.$$

**Corollary 4. [Kazdan-Warner type obstructions]** *Let  $f : S^n \rightarrow \mathbb{R}^{n+1}$ ,  $n \geq 2$ , be a conformal immersion (possibly with branching points of even order in the case  $n = 2$ ). We denote by  $H : S^n \rightarrow \mathbb{R}$  the mean curvature of this immersion. Then, for any conformal vector field  $X$  we have*

$$\int_{S^n} (\partial_X H) f^*(d\mu) = 0$$

where  $d\mu$  is the volume element on  $f(S^n)$  induced from the Euclidean metric on  $\mathbb{R}^{n+1}$ . In particular,  $\partial_X H$  changes sign.

The corollary is particularly interesting in dimension  $n = 2$ . If  $f : S^2 \rightarrow \mathbb{R}^3$  is any immersion, then after possibly composing with a diffeomorphism  $S^2 \rightarrow S^2$ , we can assume that  $f$  is conformal.

The corollary is analogous to results in [KW75], [BE87] and [DR99].

**Proof:** Let  $\psi$  be parallel spinor on  $\mathbb{R}^{n+1}$ . Then, as proven in [KS96, Ba98, Fr98], the restriction of  $\psi$  on  $\Sigma$  satisfies equation (2) with  $p = 2n/(n-1)$ , and  $|\psi|^p d\nu = f^*(d\mu)$  where  $d\nu$  is the standard volume element on  $S^n$ . Since this equation is conformally invariant we obtain a solution of (2) on  $S^n$  equipped with the standard metric. The corollary then immediately follows from example (2) above.

**Example 5.** Let  $x_3 : S^2 \rightarrow \mathbb{R}$  be the third component of the standard inclusion. One shows that  $X := \operatorname{grad} x_3$  is a conformal vector field on  $S^2$ , where the gradient is taken with respect to the standard metric on  $S^2$ . Then for any  $\varepsilon \in \mathbb{R} \setminus \{0\}$  one has  $\partial_X(\varepsilon x_3 + 1) = \varepsilon g(\operatorname{grad} x_3, \operatorname{grad} x_3)$  which is of constant sign. Hence  $\varepsilon x_3 + 1 : S^2 \rightarrow \mathbb{R}$  is not the mean curvature of a conformal immersion.

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